# Some existence results on periodic solutions of Euler-Lagrange equations in an Orlicz-Sobolev space setting 

S. Acinas ${ }^{\mathrm{a}, \mathrm{b}, *}$, L. Buri ${ }^{\text {c }}$, G. Giubergia ${ }^{\mathrm{c}}$, F. Mazzone ${ }^{\mathrm{c}, * *}$, E. Schwindt ${ }^{\mathrm{d}}$<br>${ }^{\text {a }}$ Instituto de Matemática Aplicada San Luis (CONICET-UNSL), (5700) San Luis, Argentina<br>${ }^{\mathrm{b}}$ Universidad Nacional de La Pampa, (6300) Santa Rosa, La Pampa, Argentina<br>${ }^{\text {c }}$ Dpto. de Matemática, Facultad de Ciencias Exactas, Físico-Químicas y Naturales, Universidad Nacional de Río Cuarto, (5800) Río Cuarto, Córdoba, Argentina<br>${ }^{\text {d }}$ Université d'Orléans, Laboratoire MAPMO, CNRS, UMR 7349, Fédération Denis Poisson, FR 2964, Bâtiment de Mathématiques, BP 6759, 45067 Orléans Cedex 2, France

## A R T I C L E I N F O

## Article history:

Received 23 December 2014
Accepted 4 June 2015
Communicated by Enzo Mitidieri

## Keywords:

Periodic solution
Orlicz-Sobolev spaces
Euler-Lagrange
$N$-function
Critical points


#### Abstract

In this paper we consider the problem of finding periodic solutions of certain EulerLagrange equations. We employ the direct method of the calculus of variations, i.e. we obtain solutions minimizing certain functional $I$. We give conditions which ensure that $I$ is finitely defined and differentiable on certain subsets of Orlicz-Sobolev spaces $W^{1} L^{\Phi}$ associated to an $N$-function $\Phi$. We show that, in some sense, it is necessary for the coercivity that the complementary function of $\Phi$ satisfy the $\Delta_{2^{-}}$ condition. We conclude by discussing conditions for the existence of minima of $I$.


© 2015 Elsevier Ltd. All rights reserved.

## 1. Introduction

This paper is concerned with the existence of periodic solutions of the problem

$$
\left\{\begin{array}{l}
\frac{d}{d t} D_{y} \mathcal{L}(t, \boldsymbol{u}(t), \dot{\boldsymbol{u}}(t))=D_{x} \mathcal{L}(t, \boldsymbol{u}(t), \dot{\boldsymbol{u}}(t)) \quad \text { a.e. } t \in(0, T)  \tag{1}\\
\boldsymbol{u}(0)-\boldsymbol{u}(T)=\boldsymbol{u}(0)-\boldsymbol{u}(T)=0
\end{array}\right.
$$

where $T>0, \boldsymbol{u}:[0, T] \rightarrow \mathbb{R}^{d}$ is absolutely continuous and the Lagrangian $\mathcal{L}:[0, T] \times \mathbb{R}^{d} \times \mathbb{R}^{d} \rightarrow \mathbb{R}$ is a Carathéodory and differentiable function satisfying the conditions

$$
\begin{equation*}
|\mathcal{L}(t, \boldsymbol{x}, \boldsymbol{y})| \leqslant a(|\boldsymbol{x}|)\left(b(t)+\Phi\left(\frac{|\boldsymbol{y}|}{\lambda}+f(t)\right)\right) \tag{2}
\end{equation*}
$$

[^0]\[

$$
\begin{align*}
& \left|D_{x} \mathcal{L}(t, \boldsymbol{x}, \boldsymbol{y})\right| \leqslant a(|\boldsymbol{x}|)\left(b(t)+\Phi\left(\frac{|\boldsymbol{y}|}{\lambda}+f(t)\right)\right),  \tag{3}\\
& \left|D_{y} \mathcal{L}(t, \boldsymbol{x}, \boldsymbol{y})\right| \leqslant a(|\boldsymbol{x}|)\left(c(t)+\varphi\left(\frac{|\boldsymbol{y}|}{\lambda}+f(t)\right)\right) . \tag{4}
\end{align*}
$$
\]

In these inequalities we assume that $a \in C\left(\mathbb{R}^{+}, \mathbb{R}^{+}\right), \lambda>0, \Phi$ is an $N$-function (see section Preliminaries for definitions), $\varphi$ is the right continuous derivative of $\Phi$. The non negative functions $b, c$ and $f$ satisfy that $b \in L_{1}^{1}([0, T]), c \in L_{1}^{\Psi}([0, T])$ and $f \in E_{1}^{\Phi}([0, T])$, where the Banach spaces $L_{1}^{1}([0, T]), L_{1}^{\Psi}([0, T])$ and $E_{1}^{\Phi}([0, T])$ will be defined later.

It is well known that problem (1) comes from a variational one, that is, a solution of (1) is a critical point of the action integral

$$
\begin{equation*}
I(\boldsymbol{u})=\int_{0}^{T} \mathcal{L}(t, \boldsymbol{u}(t), \dot{\boldsymbol{u}}(t)) d t \tag{5}
\end{equation*}
$$

Variational problems and Hamiltonian systems have been studied extensively. Classic references of these subjects are [23,28,14]. Problems like (1) have maintained the interest of researchers as the recent literature on the topic testifies. For Lagrangian functions of the type $\mathcal{L}(t, \boldsymbol{x}, \boldsymbol{y})=\frac{|\boldsymbol{y}|^{2}}{2}+F(t, \boldsymbol{x})$ many solvability conditions have been given expanding the results of [23]. In [29] the function $F$ was split up into two potentials, one of them with a property of subadditivity and the other with a bounded gradient. In [30] it was required a certain sublinearity condition on the gradient of the potential $F$; and, in [34] it was considered a potential $F$ given by a sum of a subconvex function and a subquadratic one. In [31] the uniform coercivity of $\int_{0}^{T} F(t, \boldsymbol{x}) d t$ was replaced by local coercivity of $F$ in some positive measure subset of $[0, T]$. In [37], the authors took a similar potential to that in [34] getting new solvability conditions and they also studied the case in which the two potentials do not have any convexity.

The Lagrangian $\mathcal{L}(t, \boldsymbol{x}, \boldsymbol{y})=\frac{|\boldsymbol{y}|^{p}}{p}+F(t, \boldsymbol{x})$ for $p>1$ was treated in more recent papers. By using the dual least action principle, in [33] it was performed the extension of some results given in [23]; and, in [32] the authors improved the work done in [34]. On the other hand, by the minimax methods in critical point theory some existence theorems were obtained. In [35] it was employed a subquadratic potential $F$ in Rabinowitz's sense and in [36] $F$ was taken as in [30].

Another source of problems, close to our proposal, is the one in which a p-Laplacian-like operator is involved. Assuming that the function $\varphi$ is a homeomorphism from $\mathbb{R}^{d}$ into itself, it is considered the differential operator $\boldsymbol{u} \mapsto\left(\varphi\left(\boldsymbol{u}^{\prime}\right)\right)^{\prime}$. In [5,6,4,22,21], using the Leray-Schauder degree theory, some existence results of solutions of equations like $\left(\varphi\left(\boldsymbol{u}^{\prime}\right)\right)^{\prime}=\boldsymbol{f}\left(t, \boldsymbol{u}(t), \boldsymbol{u}^{\prime}(t)\right)$ were obtained under different boundary conditions (periodic, Dirichlet, von Neumann) and where $\boldsymbol{f}$ is not necessarily a gradient. We point out that our approach differs from that of previous articles because we tackle the direct method of the calculus of variations.

In the Orlicz-Sobolev space setting, in [15] a constrained minimization problem associated to the existence of eigenvalues for certain differential operators involving $N$-functions was studied. Slightly away from the problems to be treated in this paper, we can mention [9,10] where A. Cianchi dealt with the regularity of minimizers of action integrals defined on several variable functions.

In this article we consider Lagrangian functions defined on Orlicz-Sobolev spaces $W^{1} L^{\Phi}$ (see $[2,19,24,25]$ ) and we use the direct method of calculus of variations. The exposition is organized as follows. In Section 2 we enumerate results related to Orlicz spaces, Orlicz-Sobolev spaces and composition operators. Almost all results in this section are essentially known. Conditions (2)-(4) are the means to ensure that $I$ is finitely defined on a non trivial subset of $W^{1} L_{d}^{\Phi}$ and $I$ is Gâteaux differentiable in this subset. We develop these issues in Theorem 3.2 of Section 3. In Section 4 we prove that critical points of (5) are solutions of (1). Conditions to guarantee the coercivity of action integrals are discussed in Section 5. Finally, our main theorem about existence of solutions of (1) is introduced and proved in Section 6.

We lay emphasis on that we use $\Delta_{2}$-condition only when necessary in a certain sense (see, for example, Lemma 5.2).

## 2. Preliminaries

For reader convenience, we give a short introduction to Orlicz and Orlicz-Sobolev spaces of vector valued functions and a list of results that we will use throughout the article. Classic references for Orlicz spaces of real valued functions are [2,19,24]. For Orlicz spaces of vector valued functions, see [27] and the references therein.

Hereafter we denote by $\mathbb{R}^{+}$the set of all non negative real numbers. A function $\Phi: \mathbb{R}^{+} \rightarrow \mathbb{R}^{+}$is called an $N$-function if $\Phi$ is given by

$$
\Phi(t)=\int_{0}^{t} \varphi(\tau) d \tau, \quad \text { for } t \geq 0
$$

where $\varphi: \mathbb{R}^{+} \rightarrow \mathbb{R}^{+}$is a right continuous non decreasing function satisfying $\varphi(0)=0, \varphi(t)>0$ for $t>0$ and $\lim _{t \rightarrow \infty} \varphi(t)=+\infty$.

Given a function $\varphi$ as above, we consider the so-called right inverse function $\psi$ of $\varphi$ which is defined by $\psi(s)=\sup _{\varphi(t) \leqslant s} t$. The function $\psi$ satisfies the same properties as the function $\varphi$, therefore we have an $N$-function $\Psi$ such that $\Psi^{\prime}=\psi$. The function $\Psi$ is called the complementary function of $\Phi$.

We say that $\Phi$ satisfies the $\Delta_{2}$-condition, denoted by $\Phi \in \Delta_{2}$, if there exist constants $K>0$ and $t_{0} \geq 0$ such that

$$
\begin{equation*}
\Phi(2 t) \leqslant K \Phi(t) \tag{6}
\end{equation*}
$$

for every $t \geq t_{0}$. If $t_{0}=0$, we say that $\Phi$ satisfies the $\Delta_{2}$-condition globally ( $\Phi \in \Delta_{2}$ globally).
Let $d$ be a positive integer. We denote by $\mathcal{M}_{d}:=\mathcal{M}_{d}([0, T])$ the set of all measurable functions defined on $[0, T]$ with values on $\mathbb{R}^{d}$ and we write $\boldsymbol{u}=\left(u_{1}, \ldots, u_{d}\right)$ for $\boldsymbol{u} \in \mathcal{M}_{d}$. In this paper we adopt the convention that bold symbols denote points in $\mathbb{R}^{d}$.

Given an $N$-function $\Phi$ we define the modular function $\rho_{\Phi}: \mathcal{M}_{d} \rightarrow \mathbb{R}^{+} \cup\{+\infty\}$ by

$$
\rho_{\Phi}(\boldsymbol{u}):=\int_{0}^{T} \Phi(|\boldsymbol{u}|) d t .
$$

Here $|\cdot|$ is the euclidean norm of $\mathbb{R}^{d}$. The Orlicz class $C_{d}^{\Phi}=C_{d}^{\Phi}([0, T])$ is given by

$$
\begin{equation*}
C_{d}^{\Phi}:=\left\{\boldsymbol{u} \in \mathcal{M}_{d} \mid \rho_{\Phi}(\boldsymbol{u})<\infty\right\} . \tag{7}
\end{equation*}
$$

The Orlicz space $L_{d}^{\Phi}=L_{d}^{\Phi}([0, T])$ is the linear hull of $C_{d}^{\Phi}$; equivalently,

$$
\begin{equation*}
L_{d}^{\Phi}:=\left\{\boldsymbol{u} \in \mathcal{M}_{d} \mid \exists \lambda>0: \rho_{\Phi}(\lambda \boldsymbol{u})<\infty\right\} . \tag{8}
\end{equation*}
$$

The Orlicz space $L_{d}^{\Phi}$ equipped with the Orlicz norm

$$
\|\boldsymbol{u}\|_{L^{\Phi}}:=\sup \left\{\int_{0}^{T} \boldsymbol{u} \cdot \boldsymbol{v} d t \mid \rho_{\Psi}(\boldsymbol{v}) \leqslant 1\right\}
$$

is a Banach space. By $\boldsymbol{u} \cdot \boldsymbol{v}$ we denote the usual dot product in $\mathbb{R}^{d}$ between $\boldsymbol{u}$ and $\boldsymbol{v}$. The following alternative expression for the norm, known as Amemiya norm, will be useful (see [19, Thm. 10.5] and [16]). For every $\boldsymbol{u} \in L^{\Phi}$,

$$
\begin{equation*}
\|\boldsymbol{u}\|_{L^{\Phi}}=\inf _{k>0} \frac{1}{k}\left\{1+\rho_{\Phi}(k \boldsymbol{u})\right\} . \tag{9}
\end{equation*}
$$

The subspace $E_{d}^{\Phi}=E_{d}^{\Phi}([0, T])$ is defined as the closure in $L_{d}^{\Phi}$ of the subspace $L_{d}^{\infty}$ of all $\mathbb{R}^{d}$-valued essentially bounded functions. It is shown that $E_{d}^{\Phi}$ is the only one maximal subspace contained in the Orlicz class $C_{d}^{\Phi}$, i.e. $\boldsymbol{u} \in E_{d}^{\Phi}$ if and only if $\rho_{\Phi}(\lambda \boldsymbol{u})<\infty$ for any $\lambda>0$.

A generalized version of Hölder's inequality holds in Orlicz spaces (see [19, Th. 9.3]). Namely, if $\boldsymbol{u} \in L_{d}^{\Phi}$ and $\boldsymbol{v} \in L_{d}^{\Psi}$ then $\boldsymbol{u} \cdot \boldsymbol{v} \in L_{1}^{1}$ and

$$
\begin{equation*}
\int_{0}^{T} \boldsymbol{v} \cdot \boldsymbol{u} d t \leqslant\|\boldsymbol{u}\|_{L^{\Phi}}\|\boldsymbol{v}\|_{L^{\psi}} \tag{10}
\end{equation*}
$$

If $X$ and $Y$ are Banach spaces such that $Y \subset X^{*}$, we denote by $\langle\cdot, \cdot\rangle: Y \times X \rightarrow \mathbb{R}$ the bilinear pairing map given by $\left\langle x^{*}, x\right\rangle=x^{*}(x)$. Hölder's inequality shows that $L_{d}^{\Psi} \subset\left[L_{d}^{\Phi}\right]^{*}$, where the pairing $\langle\boldsymbol{v}, \boldsymbol{u}\rangle$ is defined by

$$
\begin{equation*}
\langle\boldsymbol{v}, \boldsymbol{u}\rangle=\int_{0}^{T} \boldsymbol{v} \cdot \boldsymbol{u} d t \tag{11}
\end{equation*}
$$

with $\boldsymbol{u} \in L_{d}^{\Phi}$ and $\boldsymbol{v} \in L_{d}^{U}$. Unless $\Phi \in \Delta_{2}$, the relation $L_{d}^{\Psi}=\left[L_{d}^{\Phi}\right]^{*}$ will not hold. In general, it is true that $\left[E_{d}^{\Phi}\right]^{*}=L_{d}^{U}$.

Like in [19], we will consider the subset $\Pi\left(E_{d}^{\Phi}, r\right)$ of $L_{d}^{\Phi}$ given by

$$
\Pi\left(E_{d}^{\Phi}, r\right):=\left\{\boldsymbol{u} \in L_{d}^{\Phi} \mid d\left(\boldsymbol{u}, E_{d}^{\Phi}\right)<r\right\} .
$$

This set is related to the Orlicz class $C_{d}^{\Phi}$ by means of inclusions, namely,

$$
\begin{equation*}
\Pi\left(E_{d}^{\Phi}, r\right) \subset r C_{d}^{\Phi} \subset \overline{\Pi\left(E_{d}^{\Phi}, r\right)} \tag{12}
\end{equation*}
$$

for any positive $r$. If $\Phi \in \Delta_{2}$, then the sets $L_{d}^{\Phi}, E_{d}^{\Phi}, \Pi\left(E_{d}^{\Phi}, r\right)$ and $C_{d}^{\Phi}$ are equal.
We define the Sobolev-Orlicz space $W^{1} L_{d}^{\Phi}$ (see [2]) by

$$
W^{1} L_{d}^{\Phi}:=\left\{\boldsymbol{u} \mid \boldsymbol{u} \text { is absolutely continuous and } \boldsymbol{u} \in L_{d}^{\Phi}\right\} .
$$

$W^{1} L_{d}^{\Phi}$ is a Banach space when equipped with the norm

$$
\|\boldsymbol{u}\|_{W^{1} L^{\Phi}}=\|\boldsymbol{u}\|_{L^{\Phi}}+\|\dot{\boldsymbol{u}}\|_{L^{\Phi}} .
$$

For a function $\boldsymbol{u} \in L_{d}^{1}([0, T])$, we write $\boldsymbol{u}=\overline{\boldsymbol{u}}+\widetilde{\boldsymbol{u}}$ where $\overline{\boldsymbol{u}}=\frac{1}{T} \int_{0}^{T} \boldsymbol{u}(t) d t$ and $\widetilde{\boldsymbol{u}}=\boldsymbol{u}-\overline{\boldsymbol{u}}$.
As usual, if $\left(X,\|\cdot\|_{X}\right)$ is a Banach space and $\left(Y,\|\cdot\|_{Y}\right)$ is a subspace of $X$, we write $Y \hookrightarrow X$ and we say that $Y$ is embedded in $X$ when the restricted identity map $i_{Y}: Y \rightarrow X$ is bounded. That is, there exists $C>0$ such that for any $y \in Y$ we have $\|y\|_{X} \leqslant C\|y\|_{Y}$. With this notation, Hölder's inequality states that $L_{d}^{\Psi} \hookrightarrow\left[L_{d}^{\Phi}\right]^{*}$; and, it is easy to see that for every $N$-function $\Phi$ we have that $L_{d}^{\infty} \hookrightarrow L_{d}^{\Phi} \hookrightarrow L_{d}^{1}$.

Recall that a function $w: \mathbb{R}^{+} \rightarrow \mathbb{R}^{+}$is called a modulus of continuity if $w$ is a continuous increasing function which satisfies $w(0)=0$. For example, it can be easily shown that $w(s)=s \Phi^{-1}(1 / s)$ is a modulus of continuity for every $N$-function $\Phi$. We say that $\boldsymbol{u}:[0, T] \rightarrow \mathbb{R}^{d}$ has modulus of continuity $w$ when there exists a constant $C>0$ such that

$$
\begin{equation*}
|\boldsymbol{u}(t)-\boldsymbol{u}(s)| \leqslant C w(|t-s|) \tag{13}
\end{equation*}
$$

We denote by $C^{w}\left([0, T], \mathbb{R}^{d}\right)$ the space of $w$-Hölder continuous functions. This is the space of all functions satisfying (13) for some $C>0$ and it is a Banach space with norm

$$
\|\boldsymbol{u}\|_{C^{w}\left([0, T], \mathbb{R}^{d}\right)}:=\|\boldsymbol{u}\|_{L^{\infty}}+\sup _{t \neq s} \frac{|\boldsymbol{u}(t)-\boldsymbol{u}(s)|}{w(|t-s|)} .
$$

An important aspect of the theory of Sobolev spaces is related to embedding theorems. There is an extensive literature on this question in the Orlicz-Sobolev space setting, see for example [8,7,11,13,18]. The next simple lemma is essentially known and we will use it systematically. For the sake of completeness, we include a brief proof of it.

Lemma 2.1. Let $w(s):=s \Phi^{-1}(1 / s)$. Then, the following statements hold:

1. $W^{1} L^{\Phi} \hookrightarrow C^{w}\left([0, T], \mathbb{R}^{d}\right)$ and for every $\boldsymbol{u} \in W^{1} L^{\Phi}$

$$
\begin{align*}
& |\boldsymbol{u}(t)-\boldsymbol{u}(s)| \leqslant\|\dot{\boldsymbol{u}}\|_{L^{\Phi}} w(|t-s|),  \tag{14}\\
& \|\boldsymbol{u}\|_{L^{\infty}} \leqslant \Phi^{-1}\left(\frac{1}{T}\right) \max \{1, T\}\|\boldsymbol{u}\|_{W^{1} L^{\Phi}} \tag{15}
\end{align*}
$$

2. For every $\boldsymbol{u} \in W^{1} L^{\Phi}$ we have $\widetilde{\boldsymbol{u}} \in L_{d}^{\infty}$ and

$$
\begin{equation*}
\|\widetilde{\boldsymbol{u}}\|_{L^{\infty}} \leqslant T \Phi^{-1}\left(\frac{1}{T}\right)\|\dot{\boldsymbol{u}}\|_{L^{\Phi}} \quad \text { (Sobolev's inequality). } \tag{16}
\end{equation*}
$$

Proof. For $0 \leqslant s \leqslant t \leqslant T$, we get

$$
\begin{align*}
|\boldsymbol{u}(t)-\boldsymbol{u}(s)| & \leqslant \int_{s}^{t}|\dot{\boldsymbol{u}}(\tau)| d \tau \\
& \leqslant\left\|\chi_{[s, t]}\right\|_{L^{\Psi}}\|\dot{\boldsymbol{u}}\|_{L^{\Phi}} \\
& =\|\dot{\boldsymbol{u}}\|_{L^{\Phi}}(t-s) \Phi^{-1}\left(\frac{1}{t-s}\right) \tag{17}
\end{align*}
$$

using Hölder's inequality and [19, Eq. (9.11)]. This proves the inequality (14).
Since $u_{i}$ is continuous, from Mean Value Theorem for integrals, there exists $s_{i} \in[0, T]$ such that $u_{i}\left(s_{i}\right)=$ $\bar{u}_{i}$. Using this $s_{i}$ value in (14) with $u_{i}$ instead of $\boldsymbol{u}$ and taking into account that $s \Phi^{-1}(1 / s)$ is increasing, we obtain Sobolev's inequality for each $u_{i}$. The inequality (16) follows easily from the corresponding result for each component of $\boldsymbol{u}$.

On the other hand, again by Hölder's inequality and [19, Eq. (9.11)], we have

$$
\begin{equation*}
|\overline{\boldsymbol{u}}|=\frac{1}{T} \int_{0}^{T}|\boldsymbol{u}(s)| d s \leqslant \Phi^{-1}\left(\frac{1}{T}\right)\|\boldsymbol{u}\|_{L^{\Phi} .} \tag{18}
\end{equation*}
$$

From (16), (18) and the fact that $\boldsymbol{u}=\overline{\boldsymbol{u}}+\widetilde{\boldsymbol{u}}$, we obtain (15). This completes the proof of item 1.
Remark 1. As a consequence of the previous lemma, there exists a constant $C$, only dependent on $T$, such that

$$
\begin{equation*}
\|\boldsymbol{u}\|_{W^{1} L^{\Phi}} \leqslant C\left(|\overline{\boldsymbol{u}}|+\|\boldsymbol{u}\|_{L^{\Phi}}\right) \tag{19}
\end{equation*}
$$

for every $\boldsymbol{u} \in W^{1} L_{d}^{\Phi}$.
The Arzelà-Ascoli Theorem implies that $C^{w}\left([0, T], \mathbb{R}^{d}\right) \hookrightarrow C\left([0, T], \mathbb{R}^{d}\right)$ is a compact embedding (see [12, Ch. 5] for the case $w(s)=|s|^{\alpha}$ with $0<\alpha \leqslant 1$; and, if $w$ is arbitrary, the proof follows with some obvious modifications). Therefore we have the subsequent result.

Corollary 2.2. Every bounded sequence $\left\{\boldsymbol{u}_{n}\right\}$ in $W^{1} L_{d}^{\Phi}$ has an uniformly convergent subsequence.
Given a continuous function $a \in C\left(\mathbb{R}^{+}, \mathbb{R}^{+}\right)$, we define the composition operator $a: \mathcal{M}_{d} \rightarrow \mathcal{M}_{d}$ by $\boldsymbol{a}(\boldsymbol{u})(t)=a(|\boldsymbol{u}(t)|)$. We will often use the following elementary consequence of Lemma 2.1.

Corollary 2.3. If $a \in C\left(\mathbb{R}^{+}, \mathbb{R}^{+}\right)$then $\boldsymbol{a}: W^{1} L_{d}^{\Phi} \rightarrow L_{1}^{\infty}([0, T])$ is bounded. More concretely, there exists a non decreasing function $A: \mathbb{R}^{+} \rightarrow \mathbb{R}^{+}$such that $\|\boldsymbol{a}(\boldsymbol{u})\|_{L^{\infty}([0, T])} \leqslant A\left(\|\boldsymbol{u}\|_{W^{1} L^{\phi}}\right)$.

Proof. Let $\alpha \in C\left(\mathbb{R}^{+}, \mathbb{R}^{+}\right)$be a non decreasing majorant of $a$, for example $\alpha(s):=\sup _{0 \leqslant t \leqslant s} a(t)$. If $\boldsymbol{u} \in W^{1} L_{d}^{\Phi}$ then, by Lemma 2.1,

$$
a(|\boldsymbol{u}(t)|) \leqslant \alpha\left(\|\boldsymbol{u}\|_{L^{\infty}}\right) \leqslant \alpha\left(\Phi^{-1}\left(\frac{1}{T}\right) \max \{1, T\}\|\boldsymbol{u}\|_{W^{1} L^{\Phi}}\right)=: A\left(\|\boldsymbol{u}\|_{W^{1} L^{\Phi}}\right) .
$$

The next lemma is an immediate consequence of principles related to operators of Nemitskii type, see [19, Section 17].

Lemma 2.4. The composition operator $\varphi$ acts from $\Pi\left(E_{d}^{\Phi}, 1\right)$ into $C_{1}^{\Psi}$.
Proof. From [19, Lemma 9.1] we have that $\boldsymbol{\varphi}\left(B_{L^{\Phi}}(0,1)\right) \subset C_{1}^{\Psi}$, where $B_{X}\left(\boldsymbol{u}_{0}, r\right)$ is the open ball with center $\boldsymbol{u}_{0}$ and radius $r>0$ in the space $X$. Now, applying [19, Lemma 17.1], we deduce that $\boldsymbol{\varphi}$ acts from $\Pi\left(E_{d}^{\Phi}, 1\right)$ into $C_{1}^{\Psi}$.

We also need the following technical lemma.
Lemma 2.5. Let $\lambda>0$ and let $\left\{\boldsymbol{u}_{n}\right\}_{n \in \mathbb{N}}$ be a sequence of functions in $\Pi\left(E_{d}^{\Phi}, \lambda\right)$ converging to $\boldsymbol{u} \in \Pi\left(E_{d}^{\Phi}, \lambda\right)$ in the $L^{\Phi}$-norm. Then, there exist a subsequence $\boldsymbol{u}_{n_{k}}$ and a real valued function $h \in \Pi\left(E_{1}^{\Phi}([0, T]), \lambda\right)$ such that $\boldsymbol{u}_{n_{k}} \rightarrow \boldsymbol{u}$ a.e. and $\left|\boldsymbol{u}_{n_{k}}\right| \leqslant h$ a.e.

Proof. Let $r:=d\left(\boldsymbol{u}, E_{d}^{\Phi}\right), r<\lambda$. As $\boldsymbol{u}_{n}$ converges to $\boldsymbol{u}$, there exists a subsequence $\left(n_{k}\right)$ such that

$$
\left\|\boldsymbol{u}_{n_{k}}-\boldsymbol{u}\right\|_{L^{\Phi}}<\frac{\lambda-r}{2} \quad \text { and } \quad\left\|\boldsymbol{u}_{n_{k}}-\boldsymbol{u}_{n_{k+1}}\right\|_{L^{\Phi}}<2^{-(k+1)}(\lambda-r) .
$$

Let $h:[0, T] \rightarrow \mathbb{R}$ be defined by

$$
\begin{equation*}
h(x)=\left|\boldsymbol{u}_{n_{1}}(x)\right|+\sum_{k=2}^{\infty}\left|\boldsymbol{u}_{n_{k}}(x)-\boldsymbol{u}_{n_{k-1}}(x)\right| . \tag{20}
\end{equation*}
$$

As a consequence of [19, Lemma 10.1] (see [27, Thm. 5.5] for vector valued functions), we have that $d\left(\boldsymbol{v}, E_{d}^{\Phi}\right)=d\left(|\boldsymbol{v}|, E_{1}^{\Phi}\right)$ for any $\boldsymbol{v} \in L_{d}^{\Phi}$. Now

$$
d\left(\left|\boldsymbol{u}_{n_{1}}\right|, E_{1}^{\Phi}\right)=d\left(\boldsymbol{u}_{n_{1}}, E_{d}^{\Phi}\right) \leqslant d\left(\boldsymbol{u}_{n_{1}}, \boldsymbol{u}\right)+d\left(\boldsymbol{u}, E_{d}^{\Phi}\right)<\frac{\lambda+r}{2} .
$$

Then

$$
d\left(h, E_{1}^{\Phi}\right) \leqslant d\left(h,\left|\boldsymbol{u}_{n_{1}}\right|\right)+d\left(\left|\boldsymbol{u}_{n_{1}}\right|, E_{1}^{\Phi}\right)<\lambda .
$$

Therefore, $h \in \Pi\left(E_{1}^{\Phi}, \lambda\right)$ and $|h|<\infty$ a.e. We conclude that the series $\boldsymbol{u}_{n_{1}}(x)+\sum_{k=2}^{\infty}\left(\boldsymbol{u}_{n_{k}}(x)-\boldsymbol{u}_{n_{k-1}}(x)\right)$ is absolutely convergent a.e. and this fact implies that $\boldsymbol{u}_{n_{k}} \rightarrow \boldsymbol{u}$ a.e. The inequality $\left|\boldsymbol{u}_{n_{k}}\right| \leqslant h$ follows straightforwardly from the definition of $h$.

A common obstacle in Orlicz spaces, that distinguishes them from $L^{p}$ spaces, is that a sequence $\boldsymbol{u}_{n} \in L_{d}^{\Phi}$ which is uniformly bounded by $h \in L_{1}^{\Phi}$ and a.e. convergent to $\boldsymbol{u}$ is not necessarily norm convergent. Fortunately, the subspace $E_{d}^{\Phi}$ has this property.

Lemma 2.6. Suppose that $\boldsymbol{u}_{n} \in L_{d}^{\Phi}$ is a sequence such that $\boldsymbol{u}_{n} \rightarrow \boldsymbol{u}$ a.e. and assume that there exists $h \in E_{1}^{\Phi}$ with $\left|\boldsymbol{u}_{n}\right| \leqslant h$ a.e., then $\left\|\boldsymbol{u}_{n}-\boldsymbol{u}\right\|_{L^{\Phi}} \rightarrow 0$.

Proof. [24, p. 84] and [19, Thm. 10.3].
We recall some useful concepts.

Definition 2.7. Given a function $I: U \rightarrow \mathbb{R}$ where $U$ is an open set of a Banach space $X$, we say that $I$ has a Gâteaux derivative at $\boldsymbol{u} \in U$ if there exists $\boldsymbol{u}^{*} \in X^{*}$ such that for every $\boldsymbol{v} \in X$

$$
\lim _{s \rightarrow 0} \frac{I(\boldsymbol{u}+s \boldsymbol{v})-I(\boldsymbol{u})}{s}=\left\langle\boldsymbol{u}^{*}, \boldsymbol{v}\right\rangle .
$$

See [3] for details.

Definition 2.8. Let $X$ be a Banach space and let $D \subset X$. A nonlinear operator $T: D \rightarrow X^{*}$ is called demicontinuous if it is continuous when $X$ is equipped with the strong topology and $X^{*}$ with the weak* topology (see [17]).

## 3. Differentiability of action integrals in Orlicz spaces

We take a moment for discussing the relevance of the function $f$ in the inequalities (2)-(4), which are a direct generalization of the conditions [23, Eq (a), p. 10]. In particular, we are interested in seeing when for every $f \in E_{1}^{\Phi}$ there exist $b \in L_{1}^{1}$ and a constant $C>0$ such that

$$
\begin{equation*}
\Phi(s+f(t)) \leqslant C \Phi(s)+b(t) \quad \text { for every } s>0 \tag{21}
\end{equation*}
$$

If (21) is true, then we can suppose $f=0$ in the Eqs. (2) and (3) and the same considerations should be done with $\varphi(s+f(t))$.

The convexity of $\Phi$ allows us to bound $\Phi(s+f(t))$ by the expression $\frac{1}{2} \Phi(2 s)+b(t)$ where $b(t):=$ $\frac{1}{2} \Phi(2 f(t)) \in L_{1}^{1}$ and $f \in E_{1}^{\Phi}$. Therefore, we can always assume $f=0$ in (2) and (3) at the price of making the value of $\lambda$ smaller. In the special case that $\Phi \in \Delta_{2}$, the inequality (6) implies (21).

If $\Phi \notin \Delta_{2}$, then (21) may not be true. In fact, if we consider the $N$-function $\Phi(s)=e^{s}-s-1$ which does not satisfy the $\Delta_{2}$-condition and $f(t)=\ln |\ln (t)|$ for $t \in\left[0, e^{-1}\right]$, then (21) does not hold. First, note that $f(t) \geq 0$ on $\left[0, e^{-1}\right]$ and

$$
\int_{0}^{e^{-1}} \Phi(\lambda f(t)) d t \leqslant \int_{0}^{e^{-1}} e^{\lambda f(t)} d t \leqslant \int_{0}^{e^{-1}}|\ln (t)|^{\lambda} d t<\infty
$$

for $\lambda>1$, hence $f \in E_{1}^{\Phi}$. Now, suppose that there exist $b \in L_{1}^{1}$ and $C>0$ satisfying (21). From the inequality $1 / 2 e^{s} \leqslant \Phi(s)+1$, we obtain

$$
\frac{1}{2} e^{s} e^{f(t)} \leqslant \Phi(s+f(t))+1 \leqslant C e^{s}+b(t)+1 ;
$$

next, dividing by $e^{s}$ and taking $s \rightarrow \infty$, we get $\frac{1}{2}|\ln (t)| \leqslant C$ which is a contradiction.
Before addressing the main results of this section, we recall a definition.

Definition 3.1. A function $\mathcal{L}:[0, T] \times \mathbb{R}^{d} \times \mathbb{R}^{d} \rightarrow \mathbb{R}$ is a Carathéodory function if for fixed ( $\boldsymbol{x}, \boldsymbol{y}$ ) the map $t \mapsto \mathcal{L}(t, \boldsymbol{x}, \boldsymbol{y})$ is measurable and for fixed $t$ the map $(\boldsymbol{x}, \boldsymbol{y}) \mapsto \mathcal{L}(t, \boldsymbol{x}, \boldsymbol{y})$ is continuous for almost everywhere $t \in[0, T]$. We say that $\mathcal{L}(t, \boldsymbol{x}, \boldsymbol{y})$ is differentiable Carathéodory if in addition $\mathcal{L}(t, \boldsymbol{x}, \boldsymbol{y})$ is continuously differentiable with respect to $\boldsymbol{x}$ and $\boldsymbol{y}$ for almost everywhere $t \in[0, T]$.

Theorem 3.2. Let $\mathcal{L}$ be a differentiable Carathéodory function satisfying (2)-(4). Then the following statements hold:

1. The action integral given by (5) is finitely defined on $\mathcal{E}_{d}^{\Phi}(\lambda):=W^{1} L_{d}^{\Phi} \cap\left\{\boldsymbol{u} \mid \boldsymbol{u} \in \Pi\left(E_{d}^{\Phi}, \lambda\right)\right\}$.
2. The function $I$ is Gâteaux differentiable on $\mathcal{E}_{d}^{\Phi}(\lambda)$ and its derivative $I^{\prime}$ is demicontinuous from $\mathcal{E}_{d}^{\Phi}(\lambda)$ into $\left[W^{1} L_{d}^{\Phi}\right]^{*}$. Moreover, $I^{\prime}$ is given by the following expression

$$
\begin{equation*}
\left\langle I^{\prime}(\boldsymbol{u}), \boldsymbol{v}\right\rangle=\int_{0}^{T}\left\{D_{x} \mathcal{L}(t, \boldsymbol{u}, \dot{\boldsymbol{u}}) \cdot \boldsymbol{v}+D_{y} \mathcal{L}(t, \boldsymbol{u}, \dot{\boldsymbol{u}}) \cdot \dot{\boldsymbol{v}}\right\} d t . \tag{22}
\end{equation*}
$$

3. If $\Psi \in \Delta_{2}$ then $I^{\prime}$ is continuous from $\mathcal{E}_{d}^{\Phi}(\lambda)$ into $\left[W^{1} L_{d}^{\Phi}\right]^{*}$ when both spaces are equipped with the strong topology.

Proof. Let $\boldsymbol{u} \in \mathcal{E}_{d}^{\Phi}(\lambda)$. Since $\lambda \Pi\left(E_{d}^{\Phi}, r\right)=\Pi\left(E_{d}^{\Phi}, \lambda r\right)$, we have $\dot{\boldsymbol{u}} / \lambda \in \Pi\left(E_{d}^{\Phi}, 1\right)$. Thus, as $f \in E_{1}^{\Phi}$ and attending to (12), we get

$$
\begin{equation*}
|\dot{\boldsymbol{u}}| / \lambda+f \in \Pi\left(E_{1}^{\Phi}, 1\right) \subset C_{1}^{\Phi} \tag{23}
\end{equation*}
$$

By Corollary 2.3 and (2), we get

$$
|\mathcal{L}(\cdot, \boldsymbol{u}, \dot{\boldsymbol{u}})| \leqslant A\left(\|\boldsymbol{u}\|_{W^{1} L^{\Phi}}\right)\left(b+\Phi\left(\frac{|\dot{\boldsymbol{u}}|}{\lambda}+f\right)\right) \in L_{1}^{1} .
$$

This fact proves item 1.
We split up the proof of item 2 into four steps.
Step 1. The nonlinear operator $\boldsymbol{u} \mapsto D_{x} \mathcal{L}(t, \boldsymbol{u}, \dot{\boldsymbol{u}})$ is continuous from $\mathcal{E}_{d}^{\Phi}(\lambda)$ into $L_{d}^{1}([0, T])$ with the strong topology on both sets.

If $\boldsymbol{u} \in \mathcal{E}_{d}^{\Phi}(\lambda)$, from (3) and (23), we obtain

$$
\begin{equation*}
\left|D_{x} \mathcal{L}(\cdot, \boldsymbol{u}, \dot{\boldsymbol{u}})\right| \leqslant A\left(\|u\|_{W^{1} L^{\Phi}}\right)\left(b+\Phi\left(\frac{|\dot{\boldsymbol{u}}|}{\lambda}+f\right)\right) \in L_{1}^{1} . \tag{24}
\end{equation*}
$$

Let $\left\{\boldsymbol{u}_{n}\right\}_{n \in \mathbb{N}}$ be a sequence of functions in $\mathcal{E}_{d}^{\Phi}(\lambda)$ and let $\boldsymbol{u} \in \mathcal{E}_{d}^{\Phi}(\lambda)$ such that $\boldsymbol{u}_{n} \rightarrow \boldsymbol{u}$ in $W^{1} L_{d}^{\Phi}$. From $\boldsymbol{u}_{n} \rightarrow \boldsymbol{u}$ in $L_{d}^{\Phi}$, there exists a subsequence $\boldsymbol{u}_{n_{k}}$ such that $\boldsymbol{u}_{n_{k}} \rightarrow \boldsymbol{u}$ a.e.; and, as $\dot{\boldsymbol{u}}_{n} \rightarrow \boldsymbol{u} \in \mathcal{E}_{d}^{\Phi}(\lambda)$, by Lemma 2.5, there exist a subsequence of $\boldsymbol{u}_{n_{k}}$ (again denoted $\boldsymbol{u}_{n_{k}}$ ) and a function $h \in \Pi\left(E_{1}^{\Phi}, \lambda\right)$ such that $\dot{\boldsymbol{u}}_{n_{k}} \rightarrow \dot{\boldsymbol{u}}$ a.e. and $\left|\dot{\boldsymbol{u}}_{n_{k}}\right| \leqslant h$ a.e. Since $\boldsymbol{u}_{n_{k}}, k=1,2, \ldots$, is a strong convergent sequence in $W^{1} L_{d}^{\Phi}$, it is a bounded sequence in $W^{1} L_{d}^{\Phi}$. According to Lemma 2.1 and Corollary 2.3, there exists $M>0$ such that $\left\|\boldsymbol{a}\left(\boldsymbol{u}_{n_{k}}\right)\right\|_{L^{\infty}} \leqslant M, k=1,2, \ldots$. From the previous facts and (24), we get

$$
\left|D_{x} \mathcal{L}\left(\cdot, \boldsymbol{u}_{n_{k}}, \dot{u}_{n_{k}}\right)\right| \leqslant M\left(b+\Phi\left(\frac{|h|}{\lambda}+f\right)\right) \in L_{1}^{1} .
$$

On the other hand, by the continuous differentiability of $\mathcal{L}$, we have

$$
D_{x} \mathcal{L}\left(t, \boldsymbol{u}_{n_{k}}(t), \dot{u}_{n_{k}}(t)\right) \rightarrow D_{x} \mathcal{L}(t, \boldsymbol{u}(t), \dot{\boldsymbol{u}}(t)) \quad \text { for a.e. } t \in[0, T] .
$$

Applying the Dominated Convergence Theorem we conclude the proof of step 1.
Step 2. The nonlinear operator $\boldsymbol{u} \mapsto D_{y} \mathcal{L}(t, \boldsymbol{u}, \dot{\boldsymbol{u}})$ is continuous from $\mathcal{E}_{d}^{\Phi}(\lambda)$ with the strong topology into $\left[L_{d}^{\Phi}\right]^{*}$ with the weak ${ }^{*}$ topology.

Let $\boldsymbol{u} \in \mathcal{E}_{d}^{\Phi}(\lambda)$. From (23) and Lemma 2.4, it follows that

$$
\begin{equation*}
\varphi\left(\frac{|\dot{\boldsymbol{u}}|}{\lambda}+f\right) \in C_{1}^{\Psi} ; \tag{25}
\end{equation*}
$$

and, Corollary 2.3 implies $\boldsymbol{a}(\boldsymbol{u}) \in L_{1}^{\infty}$. Therefore, in virtue of (4) we get

$$
\begin{equation*}
\left|D_{y} \mathcal{L}(\cdot, \boldsymbol{u}, \dot{\boldsymbol{u}})\right| \leqslant A\left(\|\boldsymbol{u}\|_{W^{1} L^{\Phi}}\right)\left(c+\varphi\left(\frac{|\dot{\boldsymbol{u}}|}{\lambda}+f\right)\right) \in L_{1}^{\psi} . \tag{26}
\end{equation*}
$$

Note that (24), (26) and the imbeddings $W^{1} L_{d}^{\Phi} \hookrightarrow L_{d}^{\infty}$ and $L_{d}^{\Psi} \hookrightarrow\left[L_{d}^{\Phi}\right]^{*}$ imply that the second member of (22) defines an element in $\left[W^{1} L_{d}^{\Phi}\right]^{*}$.

Let $\boldsymbol{u}_{n}, \boldsymbol{u} \in \mathcal{E}_{d}^{\Phi}(\lambda)$ such that $\boldsymbol{u}_{n} \rightarrow \boldsymbol{u}$ in the norm of $W^{1} L_{d}^{\Phi}$. We must prove that $D_{y} \mathcal{L}\left(\cdot, \boldsymbol{u}_{n}, \dot{\boldsymbol{u}}_{n}\right)$ $\xrightarrow{w^{*}} D_{y} \mathcal{L}(\cdot, \boldsymbol{u}, \dot{\boldsymbol{u}})$. On the contrary, there exist $\boldsymbol{v} \in L_{d}^{\Phi}, \epsilon>0$ and a subsequence of $\left\{\boldsymbol{u}_{n}\right\}$ (denoted $\left\{\boldsymbol{u}_{n}\right\}$ for simplicity) such that

$$
\begin{equation*}
\left|\left\langle D_{y} \mathcal{L}\left(\cdot, \boldsymbol{u}_{n}, \dot{\boldsymbol{u}}_{n}\right), \boldsymbol{v}\right\rangle-\left\langle D_{y} \mathcal{L}(\cdot, \boldsymbol{u}, \dot{\boldsymbol{u}}), \boldsymbol{v}\right\rangle\right| \geq \epsilon . \tag{27}
\end{equation*}
$$

We have $\boldsymbol{u}_{n} \rightarrow \boldsymbol{u}$ in $L_{d}^{\Phi}$ and $\dot{\boldsymbol{u}}_{n} \rightarrow \dot{\boldsymbol{u}}$ in $L_{d}^{\Phi}$. By Lemma 2.5, there exist a subsequence $\boldsymbol{u}_{n_{k}}$ and a function $h \in \Pi\left(E_{1}^{\Phi}, \lambda\right)$ such that $\boldsymbol{u}_{n_{k}} \rightarrow \boldsymbol{u}$ a.e., $\dot{\boldsymbol{u}}_{n_{k}} \rightarrow \dot{\boldsymbol{u}}$ a.e. and $\left|\dot{\boldsymbol{u}}_{n_{k}}\right| \leqslant h$ a.e. As in the previous step, since $\boldsymbol{u}_{n}$ is a convergent sequence, the Corollary 2.3 implies that $a\left(\left|\boldsymbol{u}_{n}(t)\right|\right)$ is uniformly bounded by a certain constant $M>0$. Therefore, with $\boldsymbol{u}_{n_{k}}$ instead of $\boldsymbol{u}$, inequality (26) becomes

$$
\begin{equation*}
\left|D_{y} \mathcal{L}\left(\cdot, \boldsymbol{u}_{n_{k}}, \dot{u}_{n_{k}}\right)\right| \leqslant M\left(c+\varphi\left(\frac{h}{\lambda}+f\right)\right) \in L_{1}^{\Psi} \tag{28}
\end{equation*}
$$

Consequently, as $v \in L_{d}^{\Phi}$ and employing Hölder's inequality, we obtain that

$$
\sup _{k}\left|D_{y} \mathcal{L}\left(\cdot, \boldsymbol{u}_{n_{k}}, \dot{\boldsymbol{u}}_{n_{k}}\right) \cdot v\right| \in L_{1}^{1}
$$

Finally, from the Lebesgue Dominated Convergence Theorem, we deduce

$$
\begin{equation*}
\int_{0}^{T} D_{y} \mathcal{L}\left(t, \boldsymbol{u}_{n_{k}}, \dot{u}_{n_{k}}\right) \cdot \boldsymbol{v} d t \rightarrow \int_{0}^{T} D_{y} \mathcal{L}(t, \boldsymbol{u}, \dot{\boldsymbol{u}}) \cdot \boldsymbol{v} d t \tag{29}
\end{equation*}
$$

which contradicts the inequality (27). This completes the proof of step 2 .
Step 3. We will prove (22). The proof follows similar lines as [23, Thm. 1.4]. For $\boldsymbol{u} \in \mathcal{E}_{d}^{\Phi}(\lambda)$ and $\mathbf{0} \neq \boldsymbol{v} \in W^{1} L_{d}^{\Phi}$, we define the function

$$
H(s, t):=\mathcal{L}(t, \boldsymbol{u}(t)+s \boldsymbol{v}(t), \dot{\boldsymbol{u}}(t)+s \dot{\boldsymbol{v}}(t)) .
$$

From [19, Lemma 10.1] (or [27, Thm. 5.5]) we obtain that if $|\boldsymbol{u}| \leqslant|\boldsymbol{v}|$ then $d\left(\boldsymbol{u}, E_{d}^{\Phi}\right) \leqslant d\left(\boldsymbol{v}, E_{d}^{\boldsymbol{\Phi}}\right)$. Therefore, for $|s| \leqslant s_{0}:=\left(\lambda-d\left(\dot{\boldsymbol{u}}, E_{d}^{\Phi}\right)\right) /\|\boldsymbol{v}\|_{W^{1} L^{\Phi}}$ we have

$$
d\left(\dot{\boldsymbol{u}}+s \dot{\boldsymbol{v}}, E_{d}^{\Phi}\right) \leqslant d\left(|\dot{\boldsymbol{u}}|+s|\dot{\boldsymbol{v}}|, E_{1}^{\Phi}\right) \leqslant d\left(|\dot{\boldsymbol{u}}|, E_{1}^{\Phi}\right)+s\|\boldsymbol{v}\|_{L^{\Phi}}<\lambda .
$$

Thus $\dot{\boldsymbol{u}}+s \dot{\boldsymbol{v}} \in \Pi\left(E_{d}^{\Phi}, \lambda\right)$ and $|\dot{\boldsymbol{u}}|+s|\boldsymbol{v}| \in \Pi\left(E_{1}^{\Phi}, \lambda\right)$. These facts imply, in virtue of Theorem 3.2 item 1, that $I(\boldsymbol{u}+s \boldsymbol{v})$ is well defined and finite for $|s| \leqslant s_{0}$. And, using Corollary 2.3 , we also see that

$$
\|a(|\boldsymbol{u}+s \boldsymbol{v}|)\|_{L^{\infty}} \leqslant A\left(\|\boldsymbol{u}+s \boldsymbol{v}\|_{W^{1} L^{\Phi}}\right) \leqslant A\left(\|\boldsymbol{u}\|_{W^{1} L^{\Phi}}+s_{0}\|\boldsymbol{v}\|_{W^{1} L^{\Phi}}\right)=: M .
$$

Now, applying Chain Rule, (24), (26) the monotonicity of $\varphi$ and $\Phi$, the fact that $\boldsymbol{v} \in L_{d}^{\infty}$ and $\boldsymbol{v} \in L_{d}^{\Phi}$ and Hölder's inequality, we get

$$
\begin{align*}
& \left|D_{s} H(s, t)\right|=\mid D_{x} \mathcal{L}(t, \boldsymbol{u}+s \boldsymbol{v}, \dot{\boldsymbol{u}}+s \dot{\boldsymbol{v}}) \cdot \boldsymbol{v}+D_{y} \mathcal{L}(t, \boldsymbol{u}+s \boldsymbol{v}, \dot{\boldsymbol{u}}+s \dot{\boldsymbol{v}}) \cdot \dot{\boldsymbol{v} \mid} \\
& \quad \leqslant M\left[\left.\left(b(t)+\Phi\left(\frac{\mid \dot{\boldsymbol{u}\left|+s_{0}\right| \boldsymbol{v} \mid}}{\lambda}+f(t)\right)\right)|\boldsymbol{v}|+\left(c(t)+\varphi\left(\frac{|\dot{\boldsymbol{u}}|+s_{0}|\dot{\boldsymbol{v}}|}{\lambda}+f(t)\right)\right) \right\rvert\, \dot{\boldsymbol{v} \mid}\right] \in L_{1}^{1} . \tag{30}
\end{align*}
$$

Consequently, $I$ has a directional derivative and

$$
\left\langle I^{\prime}(\boldsymbol{u}), \boldsymbol{v}\right\rangle=\left.\frac{d}{d s} I(\boldsymbol{u}+s \boldsymbol{v})\right|_{s=0}=\int_{0}^{T}\left\{D_{x} \mathcal{L}(t, \boldsymbol{u}, \dot{\boldsymbol{u}}) \cdot \boldsymbol{v}+D_{y} \mathcal{L}(t, \boldsymbol{u}, \dot{\boldsymbol{u}}) \cdot \boldsymbol{v}\right\} d t .
$$

Moreover, from (24), (26), Lemma 2.1 and the previous formula, we obtain

$$
\left|\left\langle I^{\prime}(\boldsymbol{u}), \boldsymbol{v}\right\rangle\right| \leqslant\left\|D_{x} \mathcal{L}\right\|_{L^{1}}\|\boldsymbol{v}\|_{L^{\infty}}+\left\|D_{y} \mathcal{L}\right\|_{L^{\boldsymbol{w}}}\|\boldsymbol{v}\|_{L^{\Phi}} \leqslant C\|\boldsymbol{v}\|_{W^{1} L^{\Phi}}
$$

with a appropriate constant $C$. This completes the proof of the Gâteaux differentiability of $I$.
Step 4. The operator $I^{\prime}: \mathcal{E}_{d}^{\Phi}(\lambda) \rightarrow\left[W^{1} L_{d}^{\Phi}\right]^{*}$ is demicontinuous. This is a consequence of the continuity of the mappings $\boldsymbol{u} \mapsto D_{x} \mathcal{L}(t, \boldsymbol{u}, \boldsymbol{u})$ and $\boldsymbol{u} \mapsto D_{y} \mathcal{L}(t, \boldsymbol{u}, \dot{\boldsymbol{u}})$. Indeed, if $\boldsymbol{u}_{n}, \boldsymbol{u} \in \mathcal{E}_{d}^{\Phi}(\lambda)$ with $\boldsymbol{u}_{n} \rightarrow \boldsymbol{u}$ in the norm of $W^{1} L_{d}^{\Phi}$ and $\boldsymbol{v} \in W^{1} L_{d}^{\Phi}$, then

$$
\begin{aligned}
\left\langle I^{\prime}\left(\boldsymbol{u}_{n}\right), \boldsymbol{v}\right\rangle & =\int_{0}^{T}\left\{D_{x} \mathcal{L}\left(t, \boldsymbol{u}_{n}, \dot{\boldsymbol{u}}_{n}\right) \cdot \boldsymbol{v}+D_{y} \mathcal{L}\left(t, \boldsymbol{u}_{n}, \dot{\boldsymbol{u}}_{n}\right) \cdot \dot{\boldsymbol{v}}\right\} d t \\
& \rightarrow \int_{0}^{T}\left\{D_{x} \mathcal{L}(t, \boldsymbol{u}, \dot{\boldsymbol{u}}) \cdot \boldsymbol{v}+D_{y} \mathcal{L}(t, \boldsymbol{u}, \dot{\boldsymbol{u}}) \cdot \dot{\boldsymbol{v}}\right\} d t \\
& =\left\langle I^{\prime}(\boldsymbol{u}), \boldsymbol{v}\right\rangle .
\end{aligned}
$$

In order to prove item 3 , it is necessary to see that the maps $\boldsymbol{u} \mapsto D_{x} \mathcal{L}(t, \boldsymbol{u}, \dot{\boldsymbol{u}})$ and $\boldsymbol{u} \mapsto D_{y} \mathcal{L}(t, \boldsymbol{u}, \dot{\boldsymbol{u}})$ are norm continuous from $\mathcal{E}_{d}^{\Phi}(\lambda)$ into $L_{d}^{1}$ and $L_{d}^{\Psi}$ respectively. The continuity of the first map has already been proved in step 1 . Let $\boldsymbol{u}_{n}, \boldsymbol{u} \in \mathcal{E}_{d}^{\Phi}(\lambda)$ with $\left\|\boldsymbol{u}_{n}-\boldsymbol{u}\right\|_{W^{1} L^{\Phi}} \rightarrow 0$. Therefore, there exist a subsequence $\boldsymbol{u}_{n_{k}} \in \mathcal{E}_{d}^{\Phi}(\lambda)$ and a function $h \in \Pi\left(E_{1}^{\Phi}, \lambda\right)$ such that (28) holds true. And, as $\Psi \in \Delta_{2}$ then the right hand side of (28) belongs to $E_{1}^{\Psi}$. Now, invoking Lemma 2.6, we prove that from any sequence $\boldsymbol{u}_{n}$ which converges to $\boldsymbol{u}$ in $W^{1} L_{d}^{\Phi}$ we can extract a subsequence such that $D_{y} \mathcal{L}\left(t, \boldsymbol{u}_{n_{k}}, \dot{\boldsymbol{u}}_{n_{k}}\right) \rightarrow D_{y} \mathcal{L}(t, \boldsymbol{u}, \dot{\boldsymbol{u}})$ in the strong topology. The desired result is obtained by a standard argument.

The continuity of $I^{\prime}$ follows from the continuity of $D_{x} \mathcal{L}$ and $D_{y} \mathcal{L}$ using the formula (22).

## 4. Critical points and Euler-Lagrange equations

In this section we derive the Euler-Lagrange equations associated to critical points of action integrals on the subspace of $T$-periodic functions. We denote by $W^{1} L_{T}^{\Phi}$ the subspace of $W^{1} L_{d}^{\Phi}$ containing all $T$-periodic functions. As usual, when $Y$ is a subspace of the Banach space $X$, we denote by $Y^{\perp}$ the annihilator subspace of $X^{*}$, i.e. the subspace that consists of all bounded linear functions which are identically zero on $Y$.

We recall that a function $f: \mathbb{R}^{d} \rightarrow \mathbb{R}$ is called strictly convex if $f\left(\frac{\boldsymbol{x}+\boldsymbol{y}}{2}\right)<\frac{1}{2}(f(\boldsymbol{x})+f(\boldsymbol{y}))$ for $\boldsymbol{x} \neq \boldsymbol{y}$. It is well known that if $f$ is a strictly convex and differentiable function, then $D_{x} f: \mathbb{R}^{d} \rightarrow \mathbb{R}^{d}$ is a one-to-one map (see, e.g. [26, Thm. 12.17]).

Theorem 4.1. Let $\boldsymbol{u} \in \mathcal{E}_{d}^{\Phi}(\lambda)$ be a T-periodic function. The following statements are equivalent:

1. $I^{\prime}(\boldsymbol{u}) \in\left(W^{1} L_{T}^{\Phi}\right)^{\perp}$.
2. $D_{y} \mathcal{L}(t, \boldsymbol{u}(t), \dot{\boldsymbol{u}}(t))$ is an absolutely continuous function and $\boldsymbol{u}$ solves the following boundary value problem

$$
\left\{\begin{array}{l}
\frac{d}{d t} D_{y} \mathcal{L}(t, \boldsymbol{u}(t), \dot{\boldsymbol{u}}(t))=D_{x} \mathcal{L}(t, \boldsymbol{u}(t), \dot{\boldsymbol{u}}(t)) \quad \text { a.e. } t \in(0, T)  \tag{31}\\
\boldsymbol{u}(0)-\boldsymbol{u}(T)=D_{y} \mathcal{L}(0, \boldsymbol{u}(0), \dot{\boldsymbol{u}}(0))-D_{y} \mathcal{L}(T, \boldsymbol{u}(T), \dot{\boldsymbol{u}}(T))=0 .
\end{array}\right.
$$

Moreover if $D_{y} \mathcal{L}(t, x, y)$ is T-periodic with respect to the variable $t$ and strictly convex with respect to $\boldsymbol{y}$, then $D_{y} \mathcal{L}(0, \boldsymbol{u}(0), \dot{\boldsymbol{u}}(0))-D_{y} \mathcal{L}(T, \boldsymbol{u}(T), \dot{\boldsymbol{u}}(T))=0$ is equivalent to $\dot{\boldsymbol{u}}(0)=\dot{\boldsymbol{u}}(T)$.

Proof. The condition $I^{\prime}(\boldsymbol{u}) \in\left(W^{1} L_{T}^{\Phi}\right)^{\perp}$ and (22) imply

$$
\int_{0}^{T} D_{y} \mathcal{L}(t, \boldsymbol{u}(t), \dot{u}(t)) \cdot \dot{v}(t) d t=-\int_{0}^{T} D_{x} \mathcal{L}(t, \boldsymbol{u}(t), \dot{\boldsymbol{u}}(t)) \cdot \boldsymbol{v}(t) d t
$$

for every $\boldsymbol{v} \in W^{1} L_{T}^{\Phi}$. By [23, pp. 6-7] we obtain that $D_{y} \mathcal{L}(t, \boldsymbol{u}(t), \dot{\boldsymbol{u}}(t))$ is absolutely continuous and $T$ periodic, therefore it is differentiable a.e. on $[0, T]$ and the first equality of (31) holds true. This completes the proof of 1 implies 2 . The proof of 2 implies 1 follows easily from (22) and (31).

The last part of the theorem is a consequence of $D_{y} \mathcal{L}(T, \boldsymbol{u}(T), \dot{\boldsymbol{u}}(T))=D_{y} \mathcal{L}(0, \boldsymbol{u}(0), \dot{\boldsymbol{u}}(0))=D_{y} \mathcal{L}(T$, $u(T), \dot{u}(0))$ and the injectivity of $D_{y} \mathcal{L}(T, u(T), \cdot)$.

## 5. Coercivity discussion

We recall a usual definition in the context of calculus of variations.

Definition 5.1. Let $X$ be a Banach space and let $D$ be an unbounded subset of $X$. Suppose that $J: D \subset$ $X \rightarrow \mathbb{R}$. We say that $J$ is coercive if $J(u) \rightarrow+\infty$ when $\|\boldsymbol{u}\|_{X} \rightarrow+\infty$.

It is well known that coercivity is a useful ingredient in the process of establishing existence of minima. Therefore, we are interested in finding conditions which ensure the coercivity of the action integral $I$ acting on $\mathcal{E}_{d}^{\Phi}(\lambda)$. For this purpose, we need to introduce the following extra condition on the Lagrangian function $\mathcal{L}$

$$
\begin{equation*}
\mathcal{L}(t, \boldsymbol{x}, \boldsymbol{y}) \geq \alpha_{0} \Phi\left(\frac{|\boldsymbol{y}|}{\Lambda}\right)+F(t, \boldsymbol{x}), \tag{32}
\end{equation*}
$$

where $\alpha_{0}, \Lambda>0$ and $F: \mathbb{R} \times \mathbb{R}^{d} \rightarrow \mathbb{R}$ is a Carathéodory function, i.e. $F(t, \boldsymbol{x})$ is measurable with respect to $t$ for every fixed $\boldsymbol{x} \in \mathbb{R}^{d}$ and it is continuous at $\boldsymbol{x}$ for a.e. $t \in[0, T]$. We observe that, from (32) and (2), we have $F(t, \boldsymbol{x}) \leqslant a(|\boldsymbol{x}|) b_{0}(t)$ with $b_{0}(t):=b(t)+\Phi(f(t)) \in L_{1}^{1}([0, T])$. In order to guarantee that integral $\int_{0}^{T} F(t, \boldsymbol{u}) d t$ is finite for $\boldsymbol{u} \in W^{1} L^{\Phi}$, we need to assume

$$
\begin{equation*}
|F(t, \boldsymbol{x})| \leqslant a(|\boldsymbol{x}|) b_{0}(t), \quad \text { for a.e. } t \in[0, T] \text { and for every } \boldsymbol{x} \in \mathbb{R}^{d} . \tag{33}
\end{equation*}
$$

As we shall see in Theorem 5.3, when $\mathcal{L}$ satisfies (2)-(4), (32) and (33), the coercivity of the action integral $I$ is related to the coercivity of the functional

$$
\begin{equation*}
J_{C, \nu}(\boldsymbol{u}):=\rho_{\Phi}\left(\frac{\boldsymbol{u}}{\Lambda}\right)-C\|\boldsymbol{u}\|_{L^{\Phi}}^{\nu}, \tag{34}
\end{equation*}
$$

for $C, \nu>0$. If $\Phi(x)=|x|^{p} / p$ then $J_{C, \nu}$ is clearly coercive for $\nu<p$. For more general $\Phi$ the situation is more interesting as it will be shown in the following lemma.

Lemma 5.2. Let $\Phi$ and $\Psi$ be complementary $N$-functions. Then:

1. If $C \Lambda<1$, then $J_{C, 1}$ is coercive.
2. If $\Psi \in \Delta_{2}$ globally, then there exists a constant $\alpha_{\Phi}>1$ such that, for any $0<\mu<\alpha_{\Phi}$,

$$
\begin{equation*}
\lim _{\|\boldsymbol{u}\|_{L^{\Phi} \rightarrow \infty} \rightarrow \infty} \frac{\rho_{\Phi}\left(\frac{u}{\Lambda}\right)}{\|\boldsymbol{u}\|_{L^{\Phi}}^{\mu}}=+\infty \tag{35}
\end{equation*}
$$

In particular, the functional $J_{C, \mu}$ is coercive for every $C>0$ and $0<\mu<a_{\Phi}$. The constant $\alpha_{\Phi}$ is one of the so-called Matuszewska-Orlicz indices (see [20, Ch. 11]).
3. If $J_{C, 1}$ is coercive with $C \Lambda>1$, then $\Psi \in \Delta_{2}$.

Proof. By (9) we have

$$
(1-C \Lambda)\|\boldsymbol{u}\|_{L^{\Phi}}+C \Lambda\|\boldsymbol{u}\|_{L^{\Phi}}=\|\boldsymbol{u}\|_{L^{\Phi}} \leqslant \Lambda+\Lambda \rho_{\Phi}\left(\frac{\boldsymbol{u}}{\Lambda}\right),
$$

then

$$
\frac{(1-C \Lambda)}{\Lambda}\|\boldsymbol{u}\|_{L^{\Phi}}-1 \leqslant \rho_{\Phi}\left(\frac{\boldsymbol{u}}{\Lambda}\right)-C\|\boldsymbol{u}\|_{L^{\Phi}}=J_{C, 1}(\boldsymbol{u})
$$

This inequality shows that $J_{C, 1}$ is coercive and therefore item 1 is proved.
In virtue of [1, Eq. (2.8)], the $\Delta_{2}$-condition on $\Psi$, [20, Thm. 11.7] and [20, Cor. 11.6], we obtain constants $K>0$ and $\alpha_{\Phi}>1$ such that

$$
\begin{equation*}
\Phi(r s) \geq K r^{\nu} \Phi(s) \tag{36}
\end{equation*}
$$

for any $0<\nu<\alpha_{\Phi}, s \geq 0$ and $r>1$.
Let $1<\mu<\nu<\alpha_{\Phi}$ and let $r>\Lambda$ be a constant that will be specified later. Then, from (36) and (9), we get

$$
\begin{aligned}
\frac{\int_{0}^{T} \Phi\left(\frac{|\boldsymbol{u}|}{\Lambda}\right) d t}{\|\boldsymbol{u}\|_{L^{\Phi}}^{\mu}} & \geq K\left(\frac{r}{\Lambda}\right)^{\nu} \frac{\int_{0}^{T} \Phi\left(r^{-1}|\boldsymbol{u}|\right) d t}{\|\boldsymbol{u}\|_{L^{\Phi}}^{\mu}} \\
& \geq K\left(\frac{r}{\Lambda}\right)^{\nu} \frac{r^{-1}\|\boldsymbol{u}\|_{L^{\Phi}}-1}{\|\boldsymbol{u}\|_{L^{\Phi}}^{\mu}} .
\end{aligned}
$$

We choose $r=\|\boldsymbol{u}\|_{L^{\Phi}} / 2$. Since $\|\boldsymbol{u}\|_{L^{\Phi}} \rightarrow+\infty$ we can assume $\|\boldsymbol{u}\|_{L^{\Phi}}>2 \Lambda$. Thus $r>\Lambda$ and

$$
\frac{\int_{0}^{T} \Phi\left(\frac{|u|}{\Lambda}\right) d t}{\|\boldsymbol{u}\|_{L^{\Phi}}^{\mu}} \geq \frac{K}{2^{\nu} \Lambda^{\nu}}\|\boldsymbol{u}\|_{L^{\Phi}}^{\nu-\mu} \rightarrow+\infty \quad \text { as }\|\boldsymbol{u}\|_{L^{\Phi}} \rightarrow+\infty
$$

because $\nu>\mu$.
With the aim of proving item 3, we suppose that $\Psi \notin \Delta_{2}$. By [19, Thm. 4.1], there exists a sequence of real numbers $r_{n}$ such that $r_{n} \rightarrow \infty$ and

$$
\begin{equation*}
\lim _{n \rightarrow \infty} \frac{r_{n} \psi\left(r_{n}\right)}{\Psi\left(r_{n}\right)}=+\infty \tag{37}
\end{equation*}
$$

Now, we choose $\boldsymbol{u}_{n}$ such that $\left|\boldsymbol{u}_{n}\right|=\Lambda \psi\left(r_{n}\right) \chi_{\left[0, \frac{1}{W\left(r_{n}\right)}\right.}$. Then, by [19, Eq. (9.11)], we get

$$
\left\|\boldsymbol{u}_{n}\right\|_{L^{\Phi}}=\Lambda \frac{\psi\left(r_{n}\right)}{\Psi\left(r_{n}\right)} \Psi^{-1}\left(\Psi\left(r_{n}\right)\right)=\Lambda \frac{r_{n} \psi\left(r_{n}\right)}{\Psi\left(r_{n}\right)} \rightarrow \infty, \quad \text { as } n \rightarrow \infty .
$$

And, using Young's equality (see [19, Eq. (2.7)]), we have

$$
\begin{aligned}
J_{C, 1}\left(\boldsymbol{u}_{n}\right) & =\int_{0}^{T} \Phi\left(\frac{\left|\boldsymbol{u}_{n}\right|}{\Lambda}\right) d t-C\left\|\boldsymbol{u}_{n}\right\|_{L^{\Phi}} \\
& =\frac{1}{\Psi\left(r_{n}\right)}\left[\Phi\left(\psi\left(r_{n}\right)\right)-C \Lambda r_{n} \psi\left(r_{n}\right)\right] \\
& =\frac{1}{\Psi\left(r_{n}\right)}\left[r_{n} \psi\left(r_{n}\right)-\Psi\left(r_{n}\right)-C \Lambda r_{n} \psi\left(r_{n}\right)\right] \\
& =\frac{(1-C \Lambda) r_{n} \psi\left(r_{n}\right)}{\Psi\left(r_{n}\right)}-1
\end{aligned}
$$

From (37) and the condition $C \Lambda>1$, we obtain $J_{C, 1}\left(\boldsymbol{u}_{n}\right) \rightarrow-\infty$, which contradicts the coercivity of $J_{C, 1}$.

Next, we present two results that establish coercivity of action integrals under different assumptions.
Theorem 5.3. Let $\mathcal{L}$ be a Lagrangian function satisfying (2)-(4), (32) and (33). We assume the following conditions:

1. There exist a non negative function $b_{1} \in L_{1}^{1}$ and a constant $\mu>0$ such that for any $\boldsymbol{x}_{\mathbf{1}}, \boldsymbol{x}_{\mathbf{2}} \in \mathbb{R}^{d}$ and a.e. $t \in[0, T]$

$$
\begin{equation*}
\left|F\left(t, \boldsymbol{x}_{2}\right)-F\left(t, \boldsymbol{x}_{1}\right)\right| \leqslant b_{1}(t)\left(1+\left|\boldsymbol{x}_{\mathbf{2}}-\boldsymbol{x}_{\mathbf{1}}\right|^{\mu}\right) \tag{38}
\end{equation*}
$$

We suppose that $\mu<\alpha_{\Phi}$, with $\alpha_{\Phi}$ as in Lemma 5.2, in the case that $\Psi \in \Delta_{2}$; and, we suppose $\mu=1$ if $\Psi$ is an arbitrary $N$-function.
2.

$$
\begin{equation*}
\int_{0}^{T} F(t, \boldsymbol{x}) d t \rightarrow \infty \quad a s|\boldsymbol{x}| \rightarrow \infty \tag{39}
\end{equation*}
$$

3. $\Psi \in \Delta_{2}$ or, alternatively, $\alpha_{0}^{-1} T \Phi^{-1}(1 / T)\left\|b_{1}\right\|_{L^{1}} \Lambda<1$.

Then the action integral $I$ is coercive.
Proof. In the subsequent estimates, we use (32), the decomposition $\boldsymbol{u}=\overline{\boldsymbol{u}}+\tilde{\boldsymbol{u}}$, Hölder's inequality and Sobolev's inequality.

$$
\begin{align*}
I(\boldsymbol{u}) & \geq \alpha_{0} \rho_{\Phi}\left(\frac{\dot{\boldsymbol{u}}}{\Lambda}\right)+\int_{0}^{T} F(t, \boldsymbol{u}) d t \\
& =\alpha_{0} \rho_{\Phi}\left(\frac{\dot{\boldsymbol{u}}}{\Lambda}\right)+\int_{0}^{T}[F(t, \boldsymbol{u})-F(t, \overline{\boldsymbol{u}})] d t+\int_{0}^{T} F(t, \overline{\boldsymbol{u}}) d t \\
& \geq \alpha_{0} \rho_{\Phi}\left(\frac{\dot{\boldsymbol{u}}}{\Lambda}\right)-\int_{0}^{T} b_{1}(t)\left(1+|\tilde{\boldsymbol{u}}(t)|^{\mu}\right) d t+\int_{0}^{T} F(t, \overline{\boldsymbol{u}}) d t \\
& \geq \alpha_{0} \rho_{\Phi}\left(\frac{\dot{\boldsymbol{u}}}{\Lambda}\right)-\left\|b_{1}\right\|_{L^{1}}\left(1+\|\tilde{\boldsymbol{u}}\|_{L^{\infty}}^{\mu}\right)+\int_{0}^{T} F(t, \overline{\boldsymbol{u}}) d t \\
& \geq \alpha_{0} \rho_{\Phi}\left(\frac{\dot{\boldsymbol{u}}}{\Lambda}\right)-\left\|b_{1}\right\|_{L^{1}}\left(1+\left[T \Phi^{-1}\left(\frac{1}{T}\right)\right]^{\mu}\|\dot{\boldsymbol{u}}\|_{L^{\Phi}}^{\mu}\right)+\int_{0}^{T} F(t, \overline{\boldsymbol{u}}) d t \\
& =\alpha_{0} J_{C, \mu}(\dot{\boldsymbol{u}})-\left\|b_{1}\right\|_{L^{1}}+\int_{0}^{T} F(t, \overline{\boldsymbol{u}}) d t \tag{40}
\end{align*}
$$

where $C=\alpha_{0}^{-1}\left[T \Phi^{-1}(1 / T)\right]^{\mu}\left\|b_{1}\right\|_{L^{1}}$. Let $\boldsymbol{u}_{n}$ be a sequence in $\mathcal{E}_{d}^{\Phi}(\lambda)$ with $\left\|\boldsymbol{u}_{n}\right\|_{W^{1} L^{\Phi}} \rightarrow \infty$ and we have to prove that $I\left(\boldsymbol{u}_{n}\right) \rightarrow \infty$. On the contrary, suppose that for a subsequence, still denoted by $\boldsymbol{u}_{n}, I\left(\boldsymbol{u}_{n}\right)$ is upper bounded. Then, from (19) and passing to a subsequence, we can assume that $\dot{\boldsymbol{u}}_{n}$ is unbounded in $L_{d}^{\Phi}$ or $\overline{\boldsymbol{u}}_{n}$ is unbounded in $\mathbb{R}^{d}$. On the other hand, (33) and (39) imply that the integral $\int_{0}^{T} F\left(t, \overline{\boldsymbol{u}}_{n}\right) d t$ is lower bounded. These observations, the lower bound of $I$ given in (40), assumption 3 in Theorem 5.3 and Lemma 5.2 imply that $I\left(\boldsymbol{u}_{n}\right)$ is not upper bounded. This contradiction leads us to the desired result.

Based on [23] we say that $F$ satisfies the condition (A) if $F(t, \boldsymbol{x})$ is a Carathéodory function, $F$ verifies (33) and $F$ is continuously differentiable with respect to $\boldsymbol{x}$. Moreover, the next inequality holds

$$
\begin{equation*}
\left|D_{x} F(t, \boldsymbol{x})\right| \leqslant a(|\boldsymbol{x}|) b_{0}(t), \quad \text { for a.e. } t \in[0, T] \text { and for every } \boldsymbol{x} \in \mathbb{R}^{d} \tag{41}
\end{equation*}
$$

The following result was proved in [23, p. 18].
Lemma 5.4. Suppose that $F$ satisfies condition (A) and (39), $F(t, \cdot)$ is differentiable and convex a.e. $t \in[0, T]$. Then, there exists $\boldsymbol{x}_{0} \in \mathbb{R}^{d}$ such that

$$
\begin{equation*}
\int_{0}^{T} D_{x} F\left(t, x_{0}\right) d t=0 \tag{42}
\end{equation*}
$$

Theorem 5.5. Let $\mathcal{L}$ be as in Theorem 5.3 and let $F$ be as in Lemma 5.4. Moreover, assume that $\Psi \in \Delta_{2}$ or, alternatively $\alpha_{0}^{-1} T \Phi^{-1}(1 / T) a\left(\left|x_{0}\right|\right)\left\|b_{0}\right\|_{L^{1}} \Lambda<1$, with $a$ and $b_{0}$ as in (33) and $x_{0} \in \mathbb{R}^{d}$ any point satisfying (42). Then I is coercive.

Proof. Using (32), [23, Eq. (18), p. 17], the decomposition $\boldsymbol{u}=\overline{\boldsymbol{u}}+\tilde{\boldsymbol{u}}$, (42), (10) and Sobolev's inequality, we get

$$
\begin{align*}
I(\boldsymbol{u}) & \geq \alpha_{0} \rho_{\Phi}\left(\frac{\dot{\boldsymbol{u}}}{\Lambda}\right)+\int_{0}^{T} F\left(t, \boldsymbol{x}_{0}\right) d t+\int_{0}^{T} D_{x} F\left(t, \boldsymbol{x}_{0}\right) \cdot\left(\boldsymbol{u}-\boldsymbol{x}_{0}\right) d t \\
& =\alpha_{0} \rho_{\Phi}\left(\frac{\dot{\boldsymbol{u}}}{\Lambda}\right)+\int_{0}^{T} F\left(t, \boldsymbol{x}_{0}\right) d t+\int_{0}^{T} D_{x} F\left(t, \boldsymbol{x}_{0}\right) \cdot \widetilde{\boldsymbol{u}} d t+\int_{0}^{T} D_{x} F\left(t, x_{0}\right) \cdot\left(\overline{\boldsymbol{u}}-\boldsymbol{x}_{0}\right) d t \\
& =\alpha_{0} \rho_{\Phi}\left(\frac{\dot{\boldsymbol{u}}}{\Lambda}\right)+\int_{0}^{T} F\left(t, \boldsymbol{x}_{0}\right) d t+\int_{0}^{T} D_{x} F\left(t, \boldsymbol{x}_{0}\right) \cdot \widetilde{\boldsymbol{u}} d t \\
& \geq \alpha_{0} \rho_{\Phi}\left(\frac{\boldsymbol{u}}{\Lambda}\right)-a\left(\left|\boldsymbol{x}_{0}\right|\right)\left\|b_{0}\right\|_{L^{1}}-a\left(\left|\boldsymbol{x}_{0}\right|\right)\left\|b_{0}\right\|_{L^{1}} T \Phi^{-1}\left(\frac{1}{T}\right)\|\dot{\boldsymbol{u}}\|_{L^{\Phi}} \\
& =\alpha_{0} J_{C, 1}(\dot{\boldsymbol{u}})-a\left(\left|\boldsymbol{x}_{0}\right|\right)\left\|b_{0}\right\|_{L^{1}} \tag{43}
\end{align*}
$$

with $C:=\alpha_{0}^{-1} a\left(\left|x_{0}\right|\right)\left\|b_{0}\right\|_{L^{1}} T \Phi^{-1}(1 / T)$.
Let $\alpha$ be as in Corollary 2.3, i.e. $\alpha$ is a non decreasing majorant of $a$. Using that $F(t, \bar{u} / 2) \leqslant(1 / 2)$ $F(t, \boldsymbol{u})+(1 / 2) F(t,-\widetilde{\boldsymbol{u}})$ and taking into account that $\Phi$ is a non negative function, inequality (33), Hölder's inequality, Corollary 2.3 and Sobolev's inequality, we obtain

$$
\begin{align*}
I(\boldsymbol{u}) & \geq \alpha_{0} \rho_{\Phi}\left(\frac{\dot{\boldsymbol{u}}}{\Lambda}\right)+2 \int_{0}^{T} F(t, \overline{\boldsymbol{u}} / 2) d t-\int_{0}^{T} F(t,-\widetilde{\boldsymbol{u}}) d t \\
& \geq 2 \int_{0}^{T} F(t, \overline{\boldsymbol{u}} / 2) d t-\left\|b_{0}\right\|_{L^{1}}\|\boldsymbol{a}(\tilde{\boldsymbol{u}})\|_{L^{\infty}} \\
& \geq 2 \int_{0}^{T} F(t, \overline{\boldsymbol{u}} / 2) d t-\left\|b_{0}\right\|_{L^{1}} \alpha\left(\|\tilde{\boldsymbol{u}}\|_{L^{\infty}}\right) \\
& \geq 2 \int_{0}^{T} F(t, \overline{\boldsymbol{u}} / 2) d t-C_{1} \alpha\left(C_{2}\|\boldsymbol{u}\|_{L^{\Phi}}\right) \tag{44}
\end{align*}
$$

for certain constants $C_{1}, C_{2}>0$.
Finally, reasoning in a similar way to that developed in the end of the proof of Theorem 5.3 we have that $I\left(\boldsymbol{u}_{n}\right) \rightarrow \infty$.

## 6. Main result

In order to find conditions for the lower semicontinuity of $I$, we perform a little adaptation of a result of [14].

Lemma 6.1. Let $\mathcal{L}(t, \boldsymbol{x}, \boldsymbol{y})$ be a differentiable Carathéodory function. Suppose that $F$ satisfies the condition (A) and the inequality

$$
\begin{equation*}
\mathcal{L}(t, \boldsymbol{x}, \boldsymbol{y}) \geq \Phi(|\boldsymbol{y}|)+F(t, \boldsymbol{x}), \tag{45}
\end{equation*}
$$

where $\Phi$ is an $N$-function. In addition, suppose that $\mathcal{L}(t, \boldsymbol{x}, \cdot)$ is convex in $\mathbb{R}^{d}$ for each $(t, \boldsymbol{x}) \in[0, T] \times \mathbb{R}^{d}$. Let $\left\{\boldsymbol{u}_{n}\right\} \subset W^{1} L^{\Phi}$ be a sequence such that $\boldsymbol{u}_{n}$ converges uniformly to a function $\boldsymbol{u} \in W^{1} L^{\Phi}$ and $\dot{\boldsymbol{u}}_{n}$ converges in the weak topology of $L_{d}^{1}$ to $\dot{\boldsymbol{u}}$. Then

$$
\begin{equation*}
I(\boldsymbol{u}) \leqslant \liminf _{n \rightarrow \infty} I\left(\boldsymbol{u}_{n}\right) \tag{46}
\end{equation*}
$$

Proof. First, we point out that (45) and (33) imply that $I$ is defined on $W^{1} L^{\Phi}$ taking values on the interval $(-\infty,+\infty]$. Let $\left\{\boldsymbol{u}_{n}\right\}$ be a sequence satisfying the assumptions of the theorem. We define the differentiable Carathéodory function $\hat{\mathcal{L}}=\mathcal{L}-F$ and we denote by $\hat{I}$ its associated action integral. Using [14, Thm. 2.1, p. 243], we get

$$
\begin{equation*}
\int_{0}^{T} \hat{\mathcal{L}}(t, \boldsymbol{u}, \boldsymbol{u}) d t \leqslant \liminf _{n \rightarrow \infty} \int_{0}^{T} \hat{\mathcal{L}}\left(t, \boldsymbol{u}_{n}, \dot{u}_{n}\right) d t . \tag{47}
\end{equation*}
$$

Taking account of the uniform convergence of $\boldsymbol{u}_{n}$ and the fact that $F$ is a Carathéodory function, we obtain that $F\left(t, \boldsymbol{u}_{n}(t)\right) \rightarrow F(t, \boldsymbol{u}(t))$ a.e. $t \in[0, T]$. Since the sequence $\boldsymbol{u}_{n}$ is uniformly bounded, from (33) follows that there exists $g \in L_{1}^{1}([0, T])$ such that $\left|F\left(t, \boldsymbol{u}_{n}(t)\right)\right| \leqslant g(t)$. Now, by the Dominated Convergence Theorem, we have that

$$
\begin{equation*}
\lim _{n \rightarrow \infty} \int_{0}^{T} F\left(t, \boldsymbol{u}_{n}(t)\right) d t=\int_{0}^{T} F(t, \boldsymbol{u}(t)) d t . \tag{48}
\end{equation*}
$$

Finally, as a consequence of (47) and (48), we obtain (46).
Gathering our previous results we obtain existence of solutions under four different sets of assumptions. We enunciate all these alternatives in the following theorem.

Theorem 6.2. Let $\Phi$ and $\Psi$ be complementary $N$-functions. Suppose that the differentiable Carathéodory function $\mathcal{L}(t, \boldsymbol{x}, \boldsymbol{y})$ is strictly convex at $\boldsymbol{y}, D_{y} \mathcal{L}$ is $T$-periodic with respect to $T$ and (2)-(4), (32), (33) and (39) are satisfied. In addition, assume that some of the following statements hold (we recall the definitions and properties of $\alpha_{0}, b_{1}, x_{0}$ and $b_{0}$ from (32), (38), (42) and (41) respectively):

1. $\Psi \in \Delta_{2}$ and (38).
2. (38) and $\alpha_{0}^{-1} T \Phi^{-1}(1 / T)\left\|b_{1}\right\|_{L^{1}} \Lambda<1$.
3. $\Psi \in \Delta_{2}, F$ satisfies condition (A) and $F(t, \cdot)$ is convex a.e. $t \in[0, T]$.
4. As item 3 but with $\alpha_{0}^{-1} T \Phi^{-1}(1 / T) a\left(\left|x_{0}\right|\right)\left\|b_{0}\right\|_{L^{1}} \Lambda<1$ instead of $\Psi \in \Delta_{2}$.

Then, problem (1) has a solution.
Proof. First of all, note that (32), (33) and (39) imply that $I$ is, lower bounded on $W^{1} L_{T}^{\Phi}$. Let $\left\{\boldsymbol{u}_{n}\right\} \subset W^{1} L_{T}^{\Phi}$ be a minimizing sequence for the problem $\min \left\{I(\boldsymbol{u}) \mid \boldsymbol{u} \in W^{1} L_{T}^{\Phi}\right\}$. Since $I\left(\boldsymbol{u}_{n}\right), n=1,2, \ldots$ is bounded, Theorem 5.3 (or Theorem 5.5 according to which of the items $1-4$ hold true) implies that $\left\{\boldsymbol{u}_{n}\right\}$ is norm bounded in $W^{1} L_{d}^{\Phi}$. Hence, in virtue of Corollary 2.2, we can assume that $\boldsymbol{u}_{n}$ converges uniformly to a $T$-periodic continuous function $\boldsymbol{u}$. The space $L_{d}^{\Phi}$ is a predual space, concretely $L_{d}^{\Phi}=\left[E_{d}^{U}\right]^{*}$. Thus, by [24, Cor. 5, p. 148] and since $\dot{\boldsymbol{u}}_{n}$ is bounded in $L_{d}^{\Phi}$, there exists a subsequence (again denoted by $\dot{\boldsymbol{u}}_{n}$ ) such that $\dot{\boldsymbol{u}}_{n}$ converges to a function $\boldsymbol{v} \in L_{d}^{\Phi}$ in the weak* topology of $L_{d}^{\Phi}$. From this fact and the uniform convergence of $\boldsymbol{u}_{n}$ to $\boldsymbol{u}$, we obtain that

$$
\int_{0}^{T} \dot{\boldsymbol{\xi}} \cdot \boldsymbol{u} d t=\lim _{n \rightarrow \infty} \int_{0}^{T} \dot{\boldsymbol{\xi}} \cdot \boldsymbol{u}_{n} d t=-\lim _{n \rightarrow \infty} \int_{0}^{T} \boldsymbol{\xi} \cdot \dot{\boldsymbol{u}}_{n} d t=-\int_{0}^{T} \boldsymbol{\xi} \cdot \boldsymbol{v} d t
$$

for every $T$-periodic function $\boldsymbol{\xi} \in C^{\infty}\left([0, T], \mathbb{R}^{d}\right) \subset E_{d}^{\Psi}$. Thus $\boldsymbol{v}=\dot{\boldsymbol{u}}$ a.e. $t \in[0, T]$ (see [23, p. 6]) and $\boldsymbol{u} \in W^{1} L_{T}^{\Phi}$.

Now, taking into account the relations $\left[L_{d}^{1}\right]^{*}=L_{d}^{\infty} \subset E_{d}^{\psi}$ and $L_{d}^{\Phi} \subset L_{d}^{1}$, we have that $\dot{\boldsymbol{u}}_{n}$ converges to $\dot{\boldsymbol{u}}$ in the weak topology of $L_{d}^{1}$. Consequently, Lemma 6.1 applied to the $N$-function $\alpha_{0} \Phi(|\cdot| / \Lambda)$ implies that

$$
I(\boldsymbol{u}) \leqslant \liminf _{n \rightarrow \infty} I\left(\boldsymbol{u}_{n}\right)=\min _{u \in W^{1} L_{T}^{\phi}} I(\boldsymbol{u}) .
$$

Hence, $\boldsymbol{u}$ is a minimum and therefore a critical point of $I$. Finally, invoking Theorem 4.1, the proof concludes.

One of the main novelties of our work is that we obtain existence of solutions for Lagrangian functions $\mathcal{L}(t, \boldsymbol{x}, \boldsymbol{y})$ where the nonlinearity is not necessarily a power function. And, in virtue of the fact that we have not supposed that the $N$-functions $\Phi$ satisfy the $\Delta_{2}$ condition, the nonlinearity is not even bounded by power functions. For example, our main theorem can be applied to Lagrangians like

$$
\mathcal{L}(t, \boldsymbol{x}, \boldsymbol{y})=e^{|\boldsymbol{y}|}+F(t, \boldsymbol{x})
$$

when $F$ satisfies some of the alternatives 1-4 in the main theorem. In this case, considering the $N$ function $\Phi(x)=e^{x}-x-1$, conditions (2)-(4) hold. Moreover, the complementary function of $\Phi$, i.e. $\Psi(x)=(1+x) \ln (1+x)-x$, satisfies the $\Delta_{2}$-condition. We can apply our main theorem even if $\Psi \notin \Delta_{2}$, although in this case we need an additional condition on the size of certain constants.

Another contribution that we want to emphasize is the condition (38). This condition is a relaxed version of the inequality

$$
\begin{equation*}
|\nabla F(t, \boldsymbol{x})| \leqslant g(t) \tag{49}
\end{equation*}
$$

with $g \in L_{1}^{1}$, considered by Mawhin and Willem (see [23, Th. 1.5]). We point out that Mean Value Theorem for derivatives and inequality (49) imply our condition (38). In a series of papers (see for example [34,37,32]) some weaker conditions than (49) for Lagrangians of the power type, such as

$$
\begin{equation*}
\mathcal{L}(t, \boldsymbol{x}, \boldsymbol{y})=|\boldsymbol{y}|^{p}+F(t, \boldsymbol{x}), \quad p>1, \tag{50}
\end{equation*}
$$

were treated. In the aforementioned papers, inequality (49) was replaced by the weaker condition

$$
\begin{equation*}
|\nabla F(t, \boldsymbol{x})| \leqslant b(t)|x|^{\mu-1}+b_{2}(t), \quad b_{1}, b_{2} \in L^{1}, \mu \geq 1 . \tag{51}
\end{equation*}
$$

Furthermore, either $\mu$ is subcritical, that is $\mu<p$; or, assuming additional conditions on $b_{1}, \mu$ is critical, i.e. $\mu=p$. Our condition (38) is different to (51), because of this fact our results are new even for Lagrangian functions like (50). For example, let $f(x)$ be a function defined on $\mathbb{R}$, continuously differentiable with $f(x)=|x|$, for $|x|>1$; and, we define $F(t, x)=g(t)\left(f(x)+\cos \left(e^{x}\right)\right)$, where $g$ is any function in $L^{1}$. Then $F$ satisfies (38) because

$$
\left|F\left(t, x_{1}\right)-F\left(t, x_{2}\right)\right| \leqslant|g(t)|\left(K\left|x_{1}-x_{2}\right|+2\right)
$$

since $f$ is a Lipschitz function. In addition, as $f(x)=|x|$ for $|x|>1$ and $\cos \left(e^{x}\right)$ is bounded, the function $F$ satisfies the coercivity condition (39). From here, our main theorem can be applied to this function $F$. However, $\frac{d}{d x} F=g(t)\left(f^{\prime}(x)-e^{x} \sin \left(e^{x}\right)\right)$ is not bounded by any power of $x$, unless $g$ is trivial, i.e. $F$ does not satisfy (51).

Alternatively, we consider the hypothesis of convexity on $F$. We would like to emphasize that a convex function does not necessarily satisfy (38). This observation is justified by the following fact: assume that $F$ is independent of $t$ and it also satisfies (38), then $F$ is sublinear, i.e.

$$
|F(x)| \leqslant a|x|+b, \quad a>0, b>0 .
$$

In fact, inequality (38) implies that $|F(x)-F(y)| \leqslant c$ with $c>0$ provided that $|x-y| \leqslant 1$. Now, if $x \in \mathbb{R}^{n}$ we look for some $n \in \mathbb{N}$ such that $n-1 \leqslant|x|<n$. Therefore,

$$
\begin{aligned}
|F(x)| & =\left|F(x)-F\left(\frac{n-1}{n} x\right)+F\left(\frac{n-1}{n} x\right)-\cdots+F\left(\frac{1}{n} x\right)-F(0)\right| \\
& \leqslant n c \leqslant c(|x|+1) .
\end{aligned}
$$

Particularly, the function $F(t, x)=|x|^{2}$ is convex and satisfies the condition (A); nevertheless, $F$ is not sublinear and consequently it does not satisfy (38). The same function allows us to see that our condition (38) does not imply (51). And condition (38) does not imply convexity either, which is clear taking the previously studied function $F=g(t)\left(f(x)+\cos \left(e^{x}\right)\right)$.

## Acknowledgments

The first four authors are partially supported by a UNRC grant number 18/C417. The first author is partially supported by a UNSL grant number 22/F223. The third and fourth authors are partially supported by a CONICET grant number 11220090100714 CO . The fifth author is supported by a AVENTURES-ANR-12-BLAN-BS01-0001-01 and ECOS-CONICyT grant number C13E05.

## References

[1] S. Acinas, G. Giubergia, F. Mazzone, E. Schwindt, On estimates for the period of solutions of equations involving the $\phi$-Laplace operator, J. Abstr. Differ. Equ. Appl. 5 (1) (2014) 21-34.
[2] R. Adams, J. Fournier, Sobolev Spaces, Elsevier/Academic Press, Amsterdam, 2003.
[3] A. Ambrosetti, G. Prodi, A Primer of Nonlinear Analysis, Cambridge University Press, Cambridge, 1995.
[4] C. Bereanu, J. Mawhin, Periodic solutions of nonlinear perturbations of $\phi$-Laplacians with possibly bounded $\phi$, Nonlinear Anal. 68 (6) (2008) 1668-1681.
[5] C. Bereanu, J. Mawhin, Existence and multiplicity results for some nonlinear problems with singular $\phi$-Laplacian, J. Differential Equations 243 (2) (2007) 536-557.
[6] C. Bereanu, J. Mawhin, Nonhomogeneous boundary value problems for some nonlinear equations with singular $\phi$-Laplacian, J. Math. Anal. Appl. 352 (1) (2009) 218-233.
[7] A. Cianchi, Some results in the theory of Orlicz spaces and applications to variational problems, in: Nonlinear Analysis, Function Spaces and Applications, Vol. 6 (Prague, 1998), Acad. Sci. Czech. Repub., Prague, 1999, pp. 50-92.
[8] A. Cianchi, A fully anisotropic Sobolev inequality, Pacific J. Math. 196 (2) (2000) 283-295.
[9] A. Cianchi, Local boundedness of minimizers of anisotropic functionals, Ann. Inst. H. Poincaré Anal. Non Linéaire 17 (2) (2000) 147-168.
[10] A. Cianchi, N. Fusco, Gradient regularity for minimizers under general growth conditions, J. Reine Angew. Math. 507 (1999) 15-36.
[11] N. Clavero, Optimal sobolev embeddings and function spaces, 2011. http://www.maia.ub.edu/~soria/sobolev1.pdf (last accessed: 22.12.14).
[12] B. Driver, Analysis tools with applications, 2003. http://www.math.ucsd.edu/~bdriver/231-02-03/Lecture_Notes/PDE-Anal-Book/analpde1.pdf (last accessed: 22.12.14).
[13] D. Edmunds, R. Kerman, L. Pick, Optimal Sobolev imbeddings involving rearrangement-invariant quasinorms, J. Funct. Anal. 170 (2) (2000) 307-355.
[14] I. Ekeland, R. Témam, Convex Analysis and Variational Problems, Society for Industrial and Applied Mathematics (SIAM), Philadelphia, PA, 1999.
[15] M. García-Huidobro, V. Le, R. Manásevich, K. Schmitt, On principal eigenvalues for quasilinear elliptic differential operators: an Orlicz-Sobolev space setting, NoDEA Nonlinear Differential Equations Appl. 6 (2) (1999) 207-225.
[16] H. Hudzik, L. Maligranda, Amemiya norm equals Orlicz norm in general, Indag. Math. (NS) 11 (4) (2000) 573-585.
[17] T. Kato, Demicontinuity, hemicontinuity and monotonicity, Bull. Amer. Math. Soc. 70 (1964) 548-550.
[18] R. Kerman, L. Pick, Optimal Sobolev imbeddings, Forum Math. 18 (4) (2006) 535-570.
[19] M.A. Krasnosel'skiĭ, J.B. Rutickiĭ, Convex Functions and Orlicz Spaces, P. Noordhoff Ltd., Groningen, 1961.
[20] L. Maligranda, Orlicz Spaces and Interpolation, in: Seminários de Matemática (Seminars in Mathematics), vol. 5, Universidade Estadual de Campinas, Departamento de Matemática, Campinas, 1989.
[21] R. Manásevich, J. Mawhin, Periodic solutions for nonlinear systems with p-Laplacian-like operators, J. Differential Equations 145 (2) (1998) 367-393.
[22] R. Manásevich, J. Mawhin, Boundary value problems for nonlinear perturbations of vector p-Laplacian-like operators, J. Korean Math. Soc. 37 (5) (2000) 665-685.
[23] J. Mawhin, M. Willem, Critical Point Theory and Hamiltonian Systems, Springer-Verlag, New York, 1989.
[24] M.M. Rao, Z.D. Ren, Theory of Orlicz Spaces, Vol. 146, Marcel Dekker, Inc., New York, 1991.
[25] M.M. Rao, Z.D. Ren, Applications of Orlicz Spaces, Marcel Dekker, Inc., New York, 2002.
[26] R.T. Rockafellar, R. Wets, Variational Analysis, Springer-Verlag, Berlin, 1998.
[27] G. Schappacher, A notion of Orlicz spaces for vector valued functions, Appl. Math. 50 (4) (2005) 355-386.
[28] M. Struwe, Variational Methods, Springer-Verlag, Berlin, 2008.
[29] C.-L. Tang, Periodic solutions of non-autonomous second order systems with $\gamma$-quasisubadditive potential, J. Math. Anal. Appl. 189 (3) (1995) 671-675.
[30] C.-L. Tang, Periodic solutions for nonautonomous second order systems with sublinear nonlinearity, Proc. Amer. Math. Soc. 126 (11) (1998) 3263-3270.
[31] C.L. Tang, X.-P. Wu, Periodic solutions for second order systems with not uniformly coercive potential, J. Math. Anal. Appl. 259 (2) (2001) 386-397.
[32] X. Tang, X. Zhang, Periodic solutions for second-order Hamiltonian systems with a p-Laplacian, Ann. Univ. Mariae Curie-Skłodowska Sect. A 64 (1) (2010) 93-113.
[33] Y. Tian, W. Ge, Periodic solutions of non-autonomous second-order systems with a p-Laplacian, Nonlinear Anal. 66 (1) (2007) 192-203.
[34] X.-P. Wu, C.-L. Tang, Periodic solutions of a class of non-autonomous second-order systems, J. Math. Anal. Appl. 236 (2) (1999) 227-235.
[35] B. Xu, C.-L. Tang, Some existence results on periodic solutions of ordinary p-Laplacian systems, J. Math. Anal. Appl. 333 (2) (2007) 1228-1236.
[36] Y.-W. Ye, C.-L. Tang, Periodic solutions for some nonautonomous second order Hamiltonian systems, J. Math. Anal. Appl. 344 (1) (2008) 462-471.
[37] F. Zhao, X. Wu, Periodic solutions for a class of non-autonomous second order systems, J. Math. Anal. Appl. 296 (2) (2004) 422-434.


[^0]:    ${ }^{*}$ Corresponding author at: Instituto de Matemática Aplicada San Luis (CONICET-UNSL), (5700) San Luis, Argentina.
    ** Corresponding author.
    E-mail addresses: sonia.acinas@gmail.com (S. Acinas), lburi@exa.unrc.edu.ar (L. Buri), ggiubergia@exa.unrc.edu.ar (G. Giubergia), fmazzone@exa.unrc.edu.ar (F. Mazzone), leris98@gmail.com (E. Schwindt).

